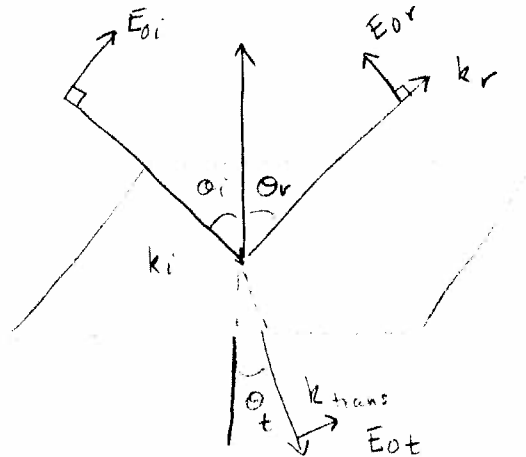




ratio of amplitudes

a) Maxwell Equations to derive r_{\perp}, r_{\parallel} ← reflected
 t_{\perp}, t_{\parallel} ← transmitted

+0.5



$$\nabla \times \vec{E} = -\omega \vec{\mu} \cdot \vec{H}(x) \quad ; \quad \nabla \times \vec{H} = -i\omega \vec{\epsilon} \cdot \vec{E}(x)$$

↑
magnetizing field
←
curl

$$\nabla \times (\nabla \times \vec{E}) = \nabla(\nabla \cdot \vec{E}) - \nabla^2 \vec{E} = \omega^2 \mu \epsilon \vec{E}(x)$$

$$\nabla \times (\nabla \times \vec{H}) = \nabla(\nabla \cdot \vec{H}) - \nabla^2 \vec{H} = \omega^2 \mu \epsilon \vec{H}(x)$$

We are assuming $\nabla \cdot \vec{B} = 0$ and $\nabla \cdot \vec{D} = \rho$
← displacement field

then $\nabla \cdot \vec{H} = \frac{1}{\mu} \nabla \cdot \vec{B}$ and $\nabla \cdot \vec{E} = \frac{1}{\epsilon} \nabla \cdot \vec{D}$

We then get the 2 diff equations

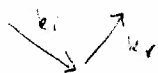
$$\nabla^2 \vec{E} + \omega^2 \epsilon \mu \vec{E} = 0 \quad \text{and} \quad \nabla^2 \vec{H} + \omega^2 \epsilon \mu \vec{H} = 0$$

The ansatz solution for the electric field is

$$\vec{E} = \vec{E}_0 e^{i(\vec{k} \cdot \vec{x} - \omega t)}$$

Where \vec{k} is the wave vector

let $|\vec{k}_i| = |\vec{k}_r|$

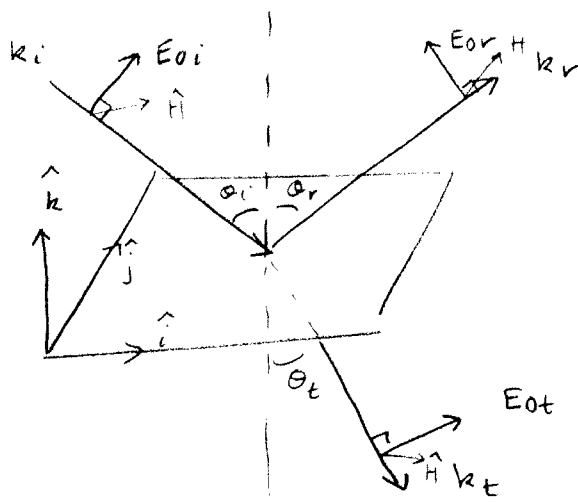


← Assuming a plane wave, no t component

Then $E_i = E_{oi} e^{i(\vec{k}_i \cdot \vec{x})}$

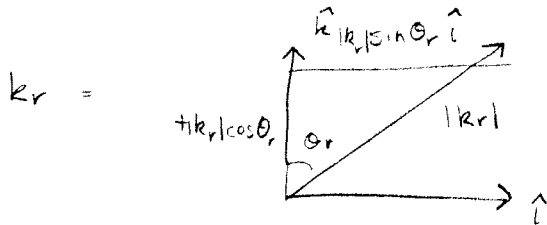
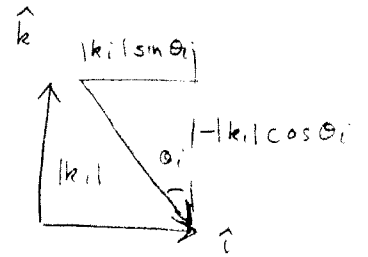
$E_r = E_{or} e^{i(\vec{k}_r \cdot \vec{x})}$

$E_t = E_{ot} e^{i(\vec{k}_t \cdot \vec{x})}$



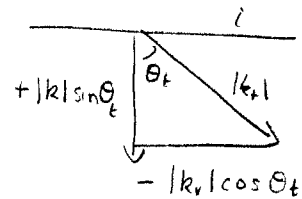
expressing \vec{k}_i in terms of \hat{i} , \hat{j} and \hat{k}

$$\vec{k}_i = |k_i| \sin \theta_i \hat{i} - |k_i| \cos \theta_i \hat{k}$$



$$\vec{k}_r = |k_r| \sin \theta_r \hat{i} + |k_r| \cos \theta_r \hat{k}$$

$$\vec{k}_t = |k_t| \sin \theta_t \hat{i} - |k_t| \cos \theta_t \hat{k}$$



$\sin k_t = \dots$

and using a similar method we can define \vec{E}_i , \vec{E}_r and \vec{E}_t in the basis set

$$\begin{aligned} E_{oi} &= |E_{oi}| \cos \theta_i \hat{i} + |E_{oi}| \sin \theta_i \hat{k} \\ E_{or} &= -|E_{or}| \cos \theta_r \hat{i} + |E_{or}| \sin \theta_r \hat{k} \\ E_{ot} &= |E_{ot}| \cos \theta_t \hat{i} + |E_{ot}| \sin \theta_t \hat{k} \end{aligned}$$

$$H_o = \frac{1}{\omega \mu} \vec{k} \times \vec{E}_o$$

$$s_o \quad H_{oi} = \frac{|k_i|}{\omega \mu_i} |E_{oi}| \hat{j}$$

$$H_{or} = -\frac{|k_r|}{\omega \mu_r} |E_{or}| \hat{j}$$

$$H_{ot} = \frac{|k_t|}{\omega \mu_t} |E_{ot}| \hat{j}$$

it must be that the magnetic field is continuous

so $H_i + H_r = H_t$

which means

$$\frac{|k_i|}{\omega \mu_i} |E_{oi}| + \frac{|k_r|}{\omega \mu_r} |E_{or}| = \frac{|k_t|}{\omega \mu_t} |E_{ot}|$$



$$n = \frac{c}{c_{\text{media}}} \quad c_{\text{media}} = \frac{\omega}{|k|}$$

so $n = \frac{c |k|}{\omega}$

$$\frac{n_i}{\mu_i} |E_{oi}| + \frac{n_r}{\mu_r} |E_{or}| = \frac{n_t}{\mu_t} |E_{ot}|$$

For continuity

$$|E_o| \cos \theta_i - |E_{or}| \cos \theta_r = |E_{ot}| \cos \theta_t$$

According to Snell's law

$$\theta_i = \theta_r \quad \text{and } n_i \mu_i = n_r \mu_r$$

and given $\mu =$

so $(|E_{oi}| - |E_{or}|) \cos \theta_i = |E_{ot}| \cos \theta_t$

$$|E_{ot}| = \frac{\mu}{n_t} \left(\frac{n_i}{\mu} |E_{oi}| + \frac{n_i}{\mu} |E_{or}| \right)$$

$$(|E_{oi}| - |E_{or}|) \cos \theta_i = \frac{\mu_t}{n_t} \left(\frac{n_i}{\mu} |E_{oi}| + \frac{n_i}{\mu} |E_{or}| \right) \cos \theta_t$$

$$\mu \left(\frac{n_t}{\mu} \cos \theta_i |E_{oi}| - \frac{n_t}{\mu} \cos \theta_i |E_{or}| \right) = \frac{n_i}{\mu} \cos \theta_t |E_{oi}| + \frac{n_i}{\mu} \cos \theta_t |E_{or}|$$

$$(n_t \cos \theta_i - n_i \cos \theta_t) |E_{oi}| = (n_i \cos \theta_t + n_t \cos \theta_i) |E_{or}|$$

$$\left(\frac{|E_{or}|}{|E_{oi}|} \right)_{||} = \frac{n_t \cos \theta_i - n_i \cos \theta_t}{n_i \cos \theta_t + n_t \cos \theta_i}$$

If we substitute for $|E_{or}|$ instead we get

$$|E_{or}| = \frac{\mu}{n_i} \left(\frac{n_t}{\mu} |E_{ot}| - \frac{n_i}{\mu} |E_{oi}| \right)$$

into

$$(|E_{oi}| - |E_{or}|) \cos \theta_i = |E_{ot}| \cos \theta_t$$

$$-|E_{or}| \cos \theta_i = |E_{ot}| \cos \theta_t - |E_{oi}| \cos \theta_i$$

$$|E_{or}| \cos \theta_i = |E_{oi}| \cos \theta_i - |E_{ot}| \cos \theta_t$$

$$\frac{1}{n_i} (n_t |E_{ot}| - n_i |E_{oi}|) \cos \theta_i = |E_{oi}| \cos \theta_i - |E_{ot}| \cos \theta_t$$

$$n_t |E_{ot}| \cos \theta_i + n_i |E_{ot}| \cos \theta_t = n_i |E_{oi}| \cos \theta_i + n_i |E_{oi}| \cos \theta_i$$

$$\left(\frac{|E_{ot}|}{|E_{oi}|} \right)_{\parallel} = \frac{2 n_i \cos \theta_i}{n_t \cos \theta_i + n_i \cos \theta_t}$$

Now looking at magnetic field \perp to the electric field

$$H_{oi} \times \hat{k} = \frac{-|k_i| |E_{oi}| \cos \theta_i}{\omega \mu} = -\frac{n_i}{\mu} |E_{oi}| \cos \theta_i$$

$$H_{or} \times \hat{k} = \frac{|k_r| |E_{or}| \cos \theta_r}{\omega \mu} = \frac{n_r}{\mu} |E_{or}| \cos \theta_r$$

$$H_{ot} \times \hat{k} = \frac{|k_t| |E_{ot}| \cos \theta_t}{\omega \mu} = \frac{n_t}{\mu} |E_{ot}| \cos \theta_t$$

$$n_i = n_r$$

$$\theta_i = \theta_r \text{ so}$$

$$\left(\frac{n_i}{\mu} (|E_{oi}| - |E_{or}|) \cos \theta_i = \frac{n_t}{\mu} |E_{ot}| \cos \theta_t \right) \mu \perp$$

$$n_i (|E_{oi}| - |E_{or}|) \cos \theta_i = n_t |E_{ot}| \cos \theta_t$$

and

$$|E_{oi}| + |E_{or}| = |E_{ot}| \quad |E_{ot}| =$$

$$n_i (|E_{oi}| - |E_{or}|) \cos \theta_i = n_t (|E_{oi}| + |E_{or}|) \cos \theta_t$$

$$n_i \cos \theta_i |E_{oi}| - n_t \cos \theta_t |E_{oi}| = n_t \cos \theta_t |E_{or}| + n_i \cos \theta_i |E_{or}|$$

$$\left(\frac{|E_{or}|}{|E_{oi}|} \right)_{\perp} = \frac{n_i \cos \theta_i - n_t \cos \theta_t}{n_i \cos \theta_i + n_t \cos \theta_t}$$

or substituting for $|E_{or}|$ instead

$$|E_{or}| = |E_{ot}| - |E_{oi}|$$

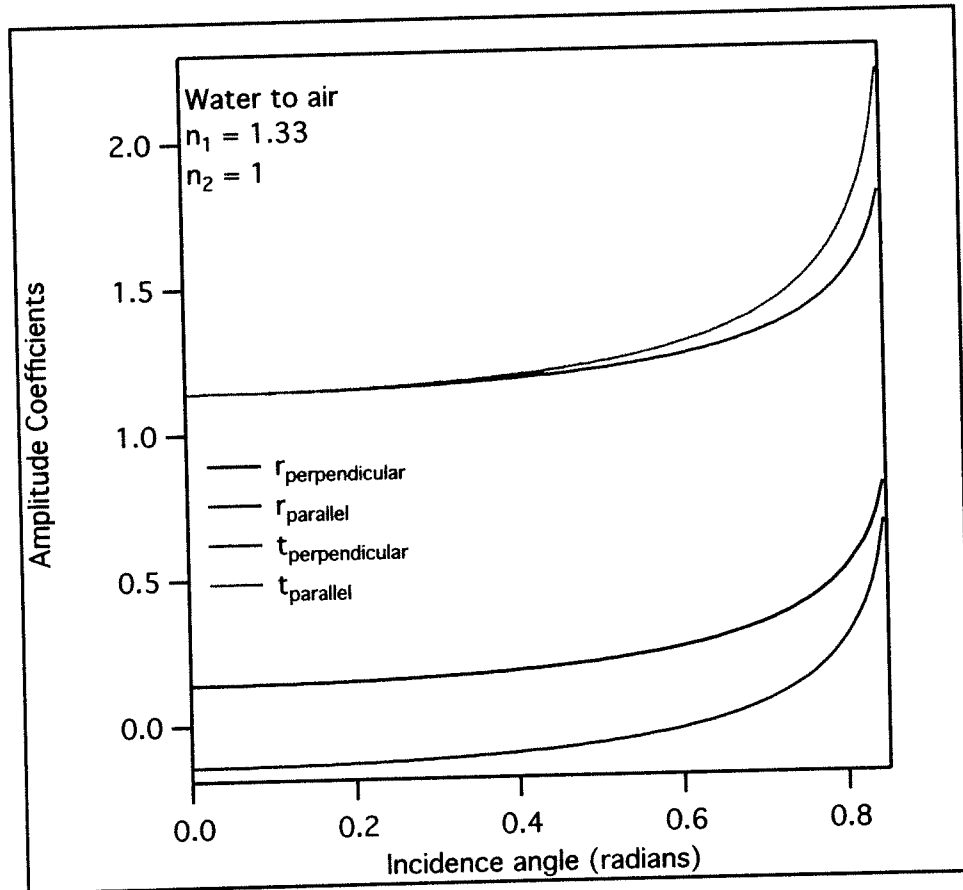
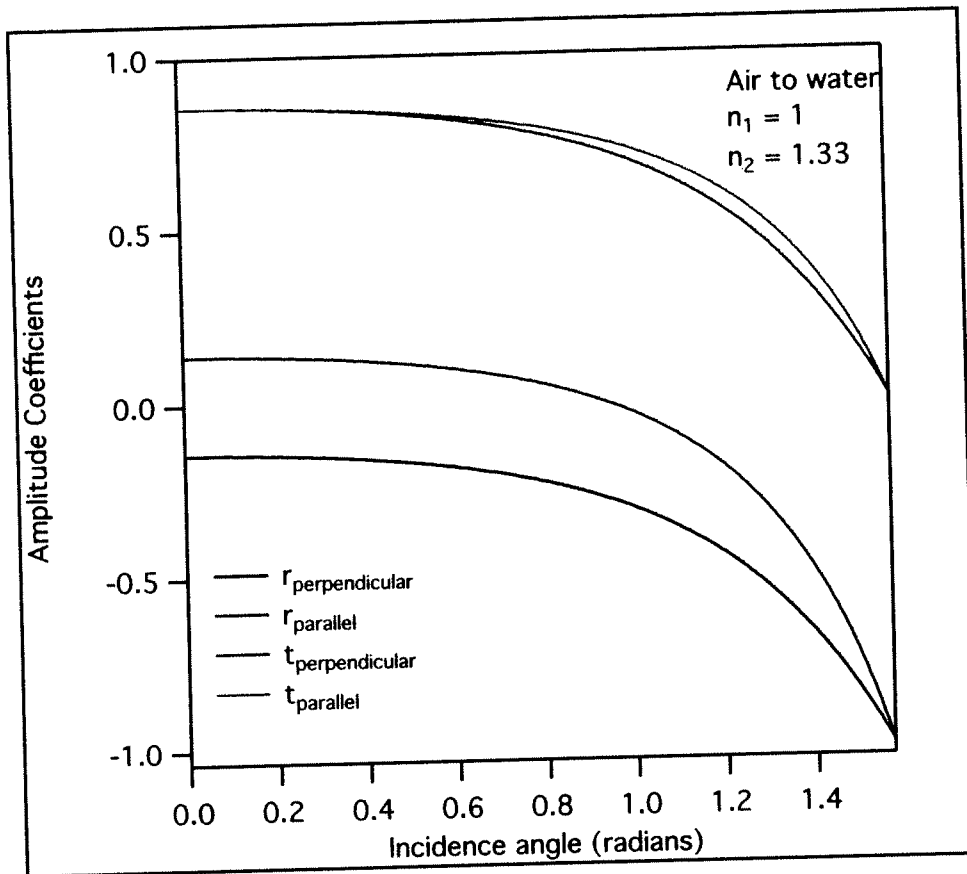
$$n_i (|E_{oi}| - |E_{or}|) \cos \theta_i = n_t \cos \theta_t |E_{ot}|$$

$$n_i \cos \theta_i |E_{or}| = n_i \cos \theta_i |E_{oi}| - n_t \cos \theta_t |E_{ot}|$$

$$n_i \cos \theta_i (|E_{ot}| - |E_{oi}|) = n_i \cos \theta_i |E_{oi}| - n_t \cos \theta_t |E_{ot}|$$

$$(n_i \cos \theta_i + n_t \cos \theta_t) |E_{ot}| = 2 n_i \cos \theta_i |E_{oi}|$$

$$\left(\frac{|E_{ot}|}{|E_{oi}|} \right)_{\perp} = \frac{2 n_i \cos \theta_i}{n_i \cos \theta_i + n_t \cos \theta_t}$$



1 c) When θ_i approaches $\sim 48.8^\circ$ (water to air), you get total internal reflection

$$n_1 \sin \theta_i = n_2 \sin \theta_t$$

$$\text{@ } \theta_t = 90^\circ, \sin \theta_t = 1$$

$$\sin \theta_i = \frac{n_2}{n_1}$$

$$\theta_i = 48.8^\circ$$

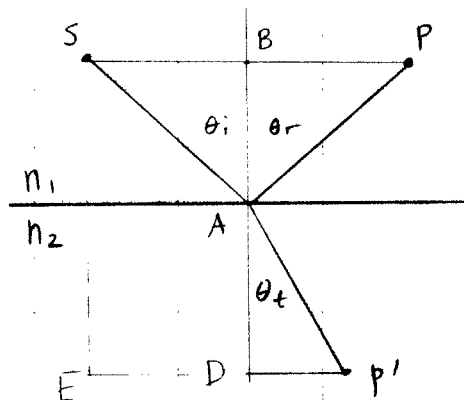
d) When swimming inside a pool, at a position just below the water surface, you can only see reflections of inside the pool & not outside.

Fermat's Principle

a)

$$\sin \theta_i = \frac{BS}{AS}$$

$$\sin \theta_t = \frac{AP'}{AP'}$$



$$c_1 = \frac{c}{n_1}$$

$$c_2 = \frac{c}{n_2}$$

$$\begin{aligned} \text{Total time, } T &= \frac{AS}{c_1} + \frac{AP'}{c_2} = \frac{\sqrt{AB^2 + BS^2}}{c_1} + \frac{\sqrt{AD^2 + DP'^2}}{c_2} \\ &= \frac{\sqrt{AB^2 + BS^2}}{c_1} + \frac{\sqrt{AD^2 + (EP' - BS)^2}}{c_2} \end{aligned}$$

$$\frac{dT}{dBS} = \frac{BS^{\sin \theta_i}}{c_1 \sqrt{AB^2 + BS^2}} + \frac{-(EP' - BS)}{c_2 \sqrt{(EP' - BS)^2 + AD^2}} \sin \theta_t = 0$$

$$\frac{d}{dx} (\sqrt{u}) = \frac{1}{2\sqrt{u}} \frac{du}{dx}$$

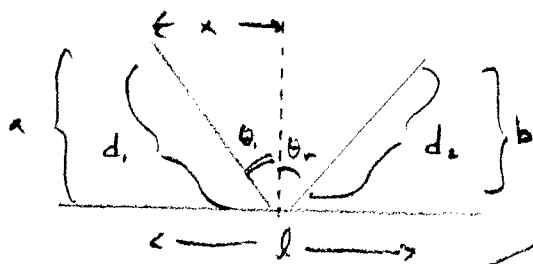
$$\frac{dT}{dBS} = \frac{\sin \theta_i}{c_1} - \frac{\sin \theta_t}{c_2} = 0 \quad \rightarrow \quad \frac{n_1}{c} \sin \theta_i = \frac{n_2}{c} \sin \theta_t$$

$$\boxed{n_1 \sin \theta_i = n_2 \sin \theta_t}$$

(b) The example uses a ray from a point source to demonstrate Snell's law. A ray is an alternative depiction of plane waves; therefore, this point source can imply Snell's law for plane waves.

(c) Test an air-to-water interface. Dip a pen into a glass of H_2O and tilt it. By observing the pen from above the H_2O , we can see that the angle of the image is distorted. This is because the indices of refraction are different in each medium, therefore $\theta_i \neq \theta_r$.

(d) $\theta_i = \theta_r$ for all cases



$$L = d_1 + d_2$$

$$= \sqrt{x^2 + a^2} + \sqrt{(l-x)^2 + b^2}$$

$$\frac{dL}{dx} = 0 = \frac{2x}{2\sqrt{x^2 + a^2}} + \frac{-2(l-x)}{2\sqrt{(l-x)^2 + b^2}}$$

minimize

$$\Rightarrow \frac{x}{\sqrt{x^2 + a^2}} = \frac{(l-x)}{\sqrt{(l-x)^2 + b^2}}$$

$$\Rightarrow \sin \theta_i = \sin \theta_r$$

$$\therefore \boxed{\theta_i = \theta_r}$$